Chapter 1

Neutron-Nucleus Interactions and Fission

The central problem in nuclear-reactor kinetics is to predict the evolution in time of the neutron population in a multiplying medium. Point kinetics allows the study of the global behaviour of the neutron population from the average properties of the medium. Before tackling, in the following chapters, the equations governing the time variation of the mactor power (proportional to the soul neutron population), we shall discuss briefly the properties of a neutron-multiplying medium. After reactions at play in a self-sustaining chain reaction, and dwell on the source of fission neutrons. Since delayed neutrons play a crucial role in mactor kinetics, we shall describe in greater detail their production in a reactor.

1.1 Neutron-Nucleus Interactions

A neutron is an elementary particle. It is stable when bound to a nucleus by the nuclear force. In its fine state, a neutron decays, with a time constant of 12 minutes, to a proton, accompanied by the emission of a β particle and an antinoutrino. In view of the fact that fire neutrons in a macro either are absorbed in matter or except, a time scale of less than 10° s, it is clear that neutron instability is completely negligible in reactor physics.

The power generated in a reactor comes from neutron-induced fissions. It results from the interaction between the neutrons and the nuclei of the physical medium. It is generally possible to neglect neutron-neutron interactions, because the density of neutrons is much smaller than that of nuclei. Free neutrons have a certain spend relative to nuclei. If one neglects the energy of thermal motion of the nuclei, the kinetic energy I of a neutron represents the energy available to a nuclear reaction between it and a nucleus. If we denote the relative spend of a neutron by v and its mass by mil, then

$$E = \frac{1}{2}mv^2 \qquad (1.1)$$

The energy is generally given in units of electron-volts (1 eV = 1.6021 x 10 $^{-8}$ J).

1.1.1 Coon Sections.

Interactions of neutrons with master are described in terms of quantities called cross sections, which can be defined in the following fashion. Imagine a unidirectional neutron beam with a density of a neutrons per cer', of speed v, impinging perpendicularly on a thin slice of a target with a density of N identical succiei per cer'. Define the neutron flux of by

$$\phi = \pi x$$
 (1-2)

The speed of neutrons in a reactor is sufficiently law that one can neglect relativistic offices. Therefore m = m_m, the neutron rest mass (1.6748 x 10⁻²⁷ kg).

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